Analysis of Liquid Flow and Mixing in an Oscillatory Flow Reactor Provided with 2D Smooth Periodic Constrictions

F. Almeida¹, F. Rocha², A. Ferreira³ *

¹Doctoral Program in Chemical and Biological Engineering, Faculty of Engineering, University of Porto, Porto, Portugal (up201202863@fe.up.pt); ²LEPABE - Laboratory for Process Engineering, Environment, Biotechnology and Energy, Department of Chemical Engineering, Faculty of Engineering, University of Porto, Porto, Portugal (frocha@fe.up.pt) ORCID 0000-0002-2778-8374; ³LEPABE - Laboratory for Process Engineering, Environment, Biotechnology and Energy, Department of Chemical Engineering, Faculty of Engineering, University of Porto, Porto, Portugal (antonio@fe.up.pt) ORCID 0000-0002-6381-981X *corresponding author

Abstract

In this research paper the residence time distribution (RTD) was monitored for a range of fluid oscillation, frequency, amplitude and flow rate in two oscillatory flow reactors (OFR) provided with 2D smooth periodic constrictions (2D-SPC) with different designs. It was studied the axial liquid dispersion using axial dispersion model and the mixing efficiency using tank-in-series model for continuous mode. Two cases, with and without fluid oscillation, were studied and determined the optimum conditions to ensure a close plug flow, an efficient mixing and a low axial liquid dispersion. The optimum operation conditions for the two 2D-SPC designs were found. Moreover, the effect of open cross-sectional area (α) on mixing and axial dispersion was also investigated. For low cross-sectional area values the mixing is higher. It was observed that fluid oscillation increases the mixing intensity even at lower flow rates, and the axial dispersion increases as the flow rate increases.

Keywords. Oscillatory Flow Reactor, Mixing, Axial Dispersion, Residence Time Distribution

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1. Introduction

In the pharmaceutical industry it is crucial to acquire an efficient mixing, as this influence the reaction yield, mass and heat transfer, as well as the quality of the final product (Ferreira, Teixeira, and Rocha 2015). However, there are still some problems related with bad mixing leading to high costs. It has been estimated that more than 600 million dollars are lost in the processes of pharmaceutical industry related with bad mixing (Paul, Atiemo-Obeng, and Kresta 2004). In order to overcome these difficulties, studies have been carried out on the development of new reactors, such as the oscillatory flow reactors (OFR). In the last years, a new generation of OFR has been arising, the meso-OFR at millimetre scale which has an ability to operate at low flow rates, reducing reagent requirements and waste (Ferreira, Teixeira, and Rocha 2015).

Several constrictions designs have been investigated, in order to achieve the optimum conditions. The first configurations were single-orifice and the multi-orifice (Mcglone et al. 2015). Since 1980s, the studies concerning the OFR increase remarkably, leading to a new constrictions designs, the smooth periodic constrictions (SPC) represented in Figure 1 (Mcglone et al. 2015).

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Figure 1: SPC configuration: Rc – Radius of curvature of the sidewall of the convergent section; Rd – Radius of curvature of the sidewall of the divergent section; L1 – Convergent and divergent section length; L2 – Straight section length; d0 – Shortest diameter of the convergent and divergent section; D – Inner diameter of the straight section (Almeida 2017)

Unlike the tubular reactors where a high superficial velocity is required to obtain a good mixing, the OFR can provide a plug flow in the laminar flow regime and operates under oscillatory flow mixing (OFM) (Phan and Harvey 2010). The oscillatory flow mixing (OFM), consists of periodically baffles inside a long tube in which either a liquid or a mixture is oscillated axially according to amplitude and frequency conditions (Brunold et al. 1989). This induces eddies on both sides of the baffle which provide both axial and radial mixing depending on the oscillatory conditions, frequency and amplitude, as shown in Figure 2.



Figure 2: Oscillatory flow mixing (Almeida 2017)

The OFM in OFR depends on several key geometrical parameters such as, the baffle spacing, l, orifice diameter, d_0 , and open cross-sectional area. The open cross-sectional area is defined as:

$$\alpha = \left(\frac{d_0}{D}\right)^2 \tag{1}$$

Phan and Harvey (2010), and Phan, Harvey, and Lavender (2011) studied the effect of new geometries (central, integral and helical) on residence time distribution. Previous studies (Dickens, Mackley, and Williams 1989; Howes and Mackley 1990; Mackley and Ni 1993), have demonstrated an improvement in residence time distribution (RTD) applying the continuous

mode. The constriction spacing, l should be in the range of 1 to 3 times the tube diameter, but l = 1.5D is the most used (Brunold et al. 1989).

Ni et al. (1998) reported that $\alpha \approx 0.25$ resulted in the shortest mixing time. Baird and Stonestreet (1995) reported a typical value of α is in the range of 0.2 to 0.4.

More recently, Ejim et al. (2017) studied the effect of the parameters α , l and D on axial dispersion in the presence of solids (polyvinyl chloride), and compared the results of axial dispersion with two constrictions, the smooth periodic constrictions (SPC) and sharp edged periodic constrictions (SEPC). Ejim et al. (2017) reported that a SPC geometry with l = 30 mm, $d_0 = 3.5$ mm, l/D = 3 and $\alpha = 0.12$ showed clear flow characteristic advantages over other meso-tubes studied.

The oscillatory flow conditions within the meso-reactor can be characterized by the three dimensionless numbers, the oscillatory Reynolds number (Re_o), Strouhal number (St) and net flow Reynolds number (Re_n), as in the previous studies.

The oscillatory Reynolds number:

$$Re_o = \frac{2\pi x_0 f d}{\nu} \tag{2}$$

At Re_o values of 100-300, eddies rings are symmetrically, but at higher Re_o the flow becomes more intensely mixed and chaotic, leading to increased axial mixing such that OFR's mixing increasingly resembles that of a stirred tank (Reis et al. 2005).

The Strouhal number:

$$St = \frac{d}{4\pi x_0} \tag{3}$$

The net flow Reynolds number:

$$Re_n = \frac{ud}{v} \tag{4}$$

The velocity ratio:

$$\varphi = \frac{Re_o}{Re_n} \tag{5}$$

Where, x_0 is the centre-to-peak amplitude (m), f is the frequency of the oscillation (Hz), d is the diameter of the tube (m), v is the kinematic viscosity (m²s⁻¹), d is the tube diameter (m), u is the superficial net flow velocity through the tube (m s⁻¹). The oscillatory Reynolds number describes the intensity of mixing applied to the tube, Strouhal number describes the effective eddy propagation inside baffle cavities and the net flow Reynolds characterizes the fluid dynamics of additional net flow within the tube (Ni and Gough 1997). It is necessary for the oscillatory flow to be dominant for the full effect of the vortex cycle to be realized. In other words, the flow should be fully reversing, so velocity ratio should be at least greater than 1 (Stonestreet and Van Der Veeken 1999).

Despite the many advantages provided by the efficient mixing of meso-OFRs, some problems related to the flow and solids suspensions have been reported. To overcome these issues, and prevent the solids deposition, a new design was developed that replaces common circular cross-section with a rectangular cross-section – Figure 4 (Ferreira et al. 2017).

In this research paper, it is explored for the first time the influence of two 2D smooth periodic constrictions (2D-SPC) designs on liquid flow and mixing. For that the residence time distribution (RTD) was monitored for a range of oscillatory conditions and flow rate.

2. Methods

2.1. Experimental apparatus

The apparatus consists of a product vessel, a meso-reactor, a feed tank, a peristaltic pump, syringe, probes, spectrometer, linear motor, light source UV-Vis and a computer.

For this study, a meso-reactor constituted by two geometries providing a total length of 2 m and a total volume of 47 ml is shown in Figure 3.



Figure 3: Schematic diagram of meso-reactor for geometry 1: 1 – Feed tank; 2 – peristaltic pump; 3 – Linear Motor; 4 – water inlet; 5 – meso-OFR-SPC; 6 – tracer injection; 7 – Probe 1; 8 - Light source UV-Vis; 9 – Spectrometer; 10 – Computer; 11 – Product vessel

The fluid oscillation was controlled by linear motor (LinMot, PS01-23x160) for a range of oscillation frequency 0 Hz to 7 Hz and amplitude 0 mm to 6 mm. The continuous net flow was controlled using a peristaltic pump (Longerpump, BT300-2J), which can feed the system with a net flow of 2.5 ml min⁻¹ ($Re_n = 9$) to 59.4 ml min⁻¹ ($Re_n = 219$). The syringe (Terumo) was used to inject a pulse tracer of indigo carmine (Merck), with a concentration of 0.2 g l⁻¹. The solutions of Indigo carmine were prepared using an analytic balance (Adam equipment, WA120). Also, experimental setup incorporates a probe where absorbance is monitored using a spectrometer (Sarspec, spec res+) and a light source UV-Vis (Sarspec). The probe was covered with an aluminum foil to reduce noise from the environment light.

In Figure 4 is shown an illustration of the 2D-SPC with a rectangular cross section.



Figure 4: 2D smooth periodic constriction: DW – tube width; d0W - shortest tube width in the constriction; ω – thickness; L1 – constriction length; L2 – Straight tube length

The 2D smooth periodic constriction characteristics are presented in Table 1.

| | D _w (mm) | dow (mm) | <i>L</i> ₁ (mm) | L2 (mm) | ω (mm) | α (%) |
|------------|------------------------|-------------|--------------------|------------|-----------|----------|
| Geometry 1 | 4 | 1.5 | 6.76 | 12.4 | 5 | 37.5 |
| Geometry 2 | 4 | 2 | 6.76 | 12.4 | 5 | 50 |

Table 1: Parameters of 2D smooth periodic constrictions The α for the 2D-smooth periodic constriction, is defined as:

$$\alpha = \frac{d_{ow}}{D_w} \tag{6}$$

Equation 1 is used just for tubes with circle cross section.

2.2. Measurement of RTD curves

In continuous mode, the meso-reactor system was constituted by two optical probes positioned at 1 m and 1.1 m downstream of the injection point of the OFR provide with geometry 1 and geometry 2, respectively. The probes were calibrated using distilled water at flow rate of 24.7 ml min⁻¹ and with desired oscillation. Before starting the experiments, it was necessary to wait one hour until the light source reaches stabilization.

After washing the meso-reactor with distilled water to eliminate air bubbles, the operational conditions of frequency, amplitude and flow rate were settled. Then at initial time it was injected a 0.2 ml of Indigo carmine (0.2 g l⁻¹) into the meso-reactor instantly, in order to achieve a perfect pulse. Hereby, the response was measured using a spectrometer and light source UV-Vis and data were recording in a computer. The data recording speed was 3 data points per second, and each data point was an average of 3 scans of the spectrometer. All the experiments were performed at room temperature (20 °C) and 3 times to achieve more precisely results. Experimental data were fitted to a plug flow with axial dispersion model and tanks-in-series model.

2.3. Absorbance calibration

For measure the absorbance over time, the Light Scan software 1.1.5 was used. The absorbance measurement was performed in wavelength of 610 nm, with an interval of \pm 25 nm, applying in the software the following expression:

$$S_1 = \int_{585}^{635} I(\lambda) \, d\lambda \tag{7}$$

Where λ is the wavelength and spectrum intensity.

Since in some cases the base line goes up, so it was necessary to find a reference value through the following expression:

$$S_2 = \int_{800}^{830} I(\lambda) \, d\lambda \tag{8}$$

For convert the absorbance into tracer concentration it was used the Lambert-Beer law:

$$A = \varepsilon lc \tag{9}$$

Where *c* is the concentration of Indigo Carmine (g l^{-1}), *l* is the light path distance (m), ε is the wavelength-dependent molar absorptivity coefficient (M⁻¹ cm⁻¹).

In Figure 5 is presented an absorbance versus concentration plot incorporated with prediction band, confidence band and the estimated line for a 99% confidence interval.

Figure 5 demonstrates a linear proportion with Indigo Carmine concentration, which shows the validity of Lambert-Beer Law for concentrations less than $0.1 \text{ g} \text{ l}^{-1}$.



The results obtained to the slope and the confidence interval at 99 % are presented in Table 2.



Figure 6: residual plots

In Figure 6 is presented the residual graphic. The plot Residuals vs Fitted of Figure 6 doesn't show a pattern. Also, the Normal Q-Q plot of Figure 6 shows that the values are closely to the line (normal distribution) so the linear model seems to be alright. The plot Scale-Location,

shows if residuals are spread equally along the ranges of predictors. As it can be seen the residuals spread wider and wider, the red smooth line isn't horizontal and shows a steep angle, so this isn't good. On the other hand, the plot Residuals vs Leverage is used to analyze the influence of outliers (points far away of linear zone) in linear regression. As it can be seen the observation 1 is outside of Cook's distance (red dashed lines). So, checking the residuals is a way to discover new insights in the model and data.

2.4. Axial dispersion model

This model has been used in RTD studies using an oscillatory flow reactor (Mackley and Ni 1993; Smith and Mackley 2006). This model was proposed by Danckwerts (1953), and it is used to describe non-ideal tubular reactors. The dispersion model in dimensionless form was described as:

$$\frac{\partial C}{\partial \theta} = \left(\frac{D_c}{uL}\right) \frac{\partial^2 C}{\partial z^2} - \frac{\partial C}{\partial z}$$
(10)

Where u is the mean axial velocity, L is the length of the test section, $\theta = ut/L$ is the dimensionless time, z = x/L is the dimensionless length, x is the axial position, t (s) is the time measured from the tracer injection, C is the dimensionless concentration, D_C (m² s⁻¹) is the axial dispersion coefficient. The dimensionless group uL/D_C is called the Peclet number (*Pe*). For a Peclet number greater than 100 occurs small deviations to a plug flow behaviour (Levenspiel 1999). The tracer concentration was normalized by:

$$C(t) = \frac{c(t)}{\int_0^\infty c(t) dt}$$
(11)

According to the close-open boundary conditions, open to diffusion to the entrance, the RTD for this model is defined as:

$$E(t) = \frac{\sqrt{\tau P e}}{2\sqrt{\pi t^3}} e^{-\frac{P e(\tau - t)^2}{4\tau t}}$$
(12)

Where $\tau = L/u$ is the mean residence time. The model parameters are Pe and τ .

2.5. Tanks-in-series model

This model assumes the concentration time response can be represented by a cascade of equal-size ideal stirred tanks, *N* (Fogler 2006). The RTD is defined as:

$$E(t) = \frac{t^{N-1}}{(N-1)!} e^{-t/\tau_i}$$
(13)

Where τ_i is the residence time in each stirred tank. The model parameters are τ_i and *N*.

3. Results and Discussion

The residence time distribution (RTD) measurements were performed at different experimental conditions. Figure 7 shows an example of the experimental data fitted with axial dispersion model. It was used the R software to fit the experimental data to the model.

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In this case the determined parameters were \tau = 23 s, Pe = 151 and N = 77.
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The experiments were performed using three types of flow conditions: Without oscillation superimposed (section 3.1), with oscillation superimposed for the two 2D-SPC designs (section 3.2), and without oscillation superimposed (section 3.3).

3.1. Without oscillation superimposed

In Figure 8 is represented the variation of D_c , P_e and N with Re_n for the two geometries.



Figure 8: a) axial dispersion coefficient, b) Number of tanks-in-series c) Number of Peclet **Figure 8** shows that the optimum conditions to ensure a plug flow behavior (high number of *Pe*), mixing intensification (high number of *N*) and a low axial dispersion coefficient are the following ones: 1) geometry 1 - $Re_n = 145$, $D_c = 2.9 \times 10^{-4}$ m² s⁻¹, N = 58, Pe = 115; 2) geometry 2 - $Re_n = 72$, $D_c = 1.8 \times 10^{-4}$ m² s⁻¹, N = 46; Pe = 91. Comparing both geometries, it has been seen that geometry 1 achieve a better mixing (N = 58) instead geometry 2 (N = 46). This may be due to the constriction in geometry 1 is tight, resulting in a high velocity in this zone, thus forming eddies in the lower pressure zones, responsible by the mixing in geometry 1. The optimum conditions assure a plug flow behavior and a mixing intensification. Zheng and Mackley (2008) observed a linear relationship of the experimental axial dispersion coefficient with net flow Reynolds number for a SPC geometry. This behavior was also observed in present work for Re_n values higher than 72 and 145 for the geometries 2 and 1, respectively.

3.2. With oscillation superimposed

Figure 9 shows the variation of D_c and P_e with Re_n and Re_o for geometry 1.



Figure 9: a) and b) x0 = 2 mm; c) and d) x0 = 4 mm; e) and f) x0 = 6 mm for geometry 1



Figure 10 shows the variation of N with Re_n and Re_o for geometry 1.

Figure 10: a) *x*0 = 2 mm; b) *x*0 = 4 mm; c) *x*0 = 6 mm for geometry 1

For geometry 1 the optimum conditions are the following ones: 1) Re_o between 200 to 400 for $x_0 = 2 \text{ mm}$; 2) Re_o between 300 to 500 for $x_0 = 4 \text{ mm}$; 3) Re_o between 400 to 600 for $x_0 = 6 \text{ mm}$. Zheng and Mackley (2008) obtained optimum conditions for $x_0 = 1 \text{ mm}$, Re_o between 100 and 300. Dickens, Mackley, and Williams (1989) obtained similar results.

It is further noted that D_c and N tend to increase with increasing of the flow rate because as fluid passes through the constriction increases the velocity causing a greater axial dispersion, also enhancing the formation of eddies. However, for $Re_n = 145$ and $x_0 = 2$ mm doesn't seem to show that. This may be due to an alteration to the flow regime or it may be a critical value, in other words, from this flow rate value can form eddies that ensure the reduction of the axial dispersion, and mixing intensification becomes high.

-Ren=9 Ren=9 1,2E-03 200 Ren=19 Ren=19 180 Ren=58 1,0E-03 Ren=58 160 Ren=72 Ren=72 a¹⁴⁰ 8,0E-04 - Ren=145 Ren=145 $D_{C} (m^{2}s^{-1})$ 120 Ren=219 Ren=219 6,0E-04 100 80 4,0E-04 60 40 2,0E-04 20 0,0E+00 0 0 200 400 600 0 100 200 300 400 500 $\operatorname{Re}_{\mathrm{o}}$ Re_{o} b) a) 1,4E-03 250 Ren=19 Ren=19 Ren=9 1,2E-03 Ren=9 200 Ren=58 Ren=58 1,0E-03 Ren=72 Ren=72 Ren=145 $D_{C} (m^{2}s^{-1})$ Ren=145 150 8,0E-04 Ren=219 Ren=219 Pe 6,0E-04 100 4,0E-04 50 2,0E-04 0,0E+00 0 0 500 1000 0 500 1000 Re_{o} Re_o d) c) 1,2E-03 250 Ren=9 Ren=9 Ren=19 - Ren=19 1,0E-03 Ren=58 - Ren=58 200 Ren=72 Ren=72 8,0E-04 Ren=145 $D_{c} (m^{2}s^{-1})$ Ren=145 150 Ren=219 - Ren=219 6,0E-04 Pe 100 4,0E-04 50 2,0E-04 0,0E+00 0 0 500 1000 1500 0 1000 500 1500 Re_{o} Re_{o} f) e)

Figure 11 demonstrates the variation of D_c and P_e with Re_n and Re_o for geometry 2.

Figure 11: a) and b) x0 = 2 mm; c) and d) x0 = 4 mm; e) and f) x0 = 6 mm for geometry 2



In Figure 12 is represented the variation of N with Re_n and Re_o for geometry 2.

Figure 12: a) *x*0 = 2 mm; b) *x*0 = 4 mm; c) *x*0 = 6 mm for geometry 2

From the analysis of Figure 11 and Figure 12 the optimum conditions are the following ones: 1) $x_0 = 2 \text{ mm}$, Re_0 between 50 and 170; 2) $x_0 = 4 \text{ mm}$, Re_0 between 530 and 780; 3) $x_0 = 6 \text{ mm}$, Re_0 above 1150. Globally, the best operating condition in this OFR-2D-SPC is Re_n de 58 or 72, $f = 7 \text{ Hz} = x_0 = 6 \text{ mm}$. It is noticed that geometry 2 presents a plug flow behavior and an improvement in mixing intensification for lower flow rates. Therefore, the geometry 2 has a positive effect for lower flow rates, because when a low flow rate fluid crosses the constriction causes a lower velocity in axial direction, since d_{0w} is largest than geometry 1, consequently, the axial dispersion coefficient is reduced.

For two 2D-SPC geometries, when an oscillation frequency is introduced, for the same Re_n and x_0 the axial dispersion coefficient tends to decrease to a minimum value. This decrease occurs because with the increase of the frequency increases the radial dispersion and a formation and faster destruction of the eddies becomes more pronounced, making the flow regime more chaotic. The optimum conditions assure a plug flow behavior and a mixing intensification for both 2D-SPC designs tested.

3.3. Without oscillation superimposed for lower flow rate

Figure 13 shows the RTD results obtained with oscillation and without oscillation of the fluid for geometry 2.



Figure 13: a) the RTD for Ren = 19 with oscillation - x_0 = 2 mm e f = 1 Hz for geometry 2; b) the RTD for Ren = 19 without oscillation for geometry 2

The analysis of the RTD profile from Figure13b, at low flow rates ($Re_n = 9$ and $Re_n = 19$) the reactor behaves as a laminar flow with low mixing, with low eddy formation. For geometry 2 with no oscillation and at low flow rates ($Re_n = 9$ and $Re_n = 19$) the results for the RTD profile were the same as Figure13b.

Therefore, it can be said that the oscillation of the fluid plays an important role in the mixing intensification of the fluid.

4. Conclusions

A continuous oscillatory flow reactor provided with 2D smooth periodic constrictions was developed for potential application in pharmaceutical industry. The performance of the RTD is an important parameter in the characterization of the OFR, since the RTD reflects the behavior of the mixing that occurs inside the OFR.

It was noticed a mixing intensification even at low flow rate, when oscillation of the fluid was present. The frequency has a radial dispersion effect, while amplitude has an axial dispersion effect. Moreover, it was concluded a positive effect for geometry 2, since it was achieved better results than geometry 1 at low flow rates when fluid oscillation is present.

Furthermore, it was concluded that, globally, the optimum conditions that assure a plug flow behavior, an efficient mixing and a low axial dispersion are as follows: 1) geometry 1 - Re_n = 19, x_0 = 2 mm and f = 5 Hz; 2) geometry 2 - Re_n = 58 or 72, x_0 = 6 mm and f = 7 Hz.

The oscillatory flow reactor provided with 2D - smooth periodic constrictions demonstrates potential in increasing process efficiency, mixing intensification and reduces the possibility of formation of secondary products, which gives it application in products of high added value in several areas, specifically in the fine chemical, biological and pharmaceutical industry.

Future works need to be performed in order to verify the 2D - smooth periodic constrictions influence on solids mixing and suspension at different scales.

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